

Today 2/7

- 1) look at characteristic curves
- 2) Derive a conservation
- 3) Introduce a Finite Volume  
"framework"

$$\text{Take } \frac{\partial u}{\partial t} + \frac{\partial (f(u))}{\partial x} = 0$$

$$\Rightarrow \frac{\partial f}{\partial t} + f'(u) \frac{\partial u}{\partial x} = 0$$

Let  $x(t)$  be a curve such that

$$\frac{dx(t)}{dt} = f'(u)$$

$$\begin{aligned} \text{Then } \frac{du(x(t), t)}{dt} &= \frac{\partial u}{\partial t} + \frac{\partial u}{\partial x} \frac{dx}{dt} \\ &= u_t + f'(u) u_x \\ &= 0 \end{aligned}$$

$$\Rightarrow \left\{ \begin{array}{l} \frac{du}{dt} = 0 \\ \frac{dx}{dt} = 0 \end{array} \right. \Rightarrow$$

$$u = \text{const}$$

$$x = ut + d$$

$\Rightarrow$  straight lines

each curve is

$$x(t) = u(x_0, 0) \cdot t + x_0$$

from each  $x_0$

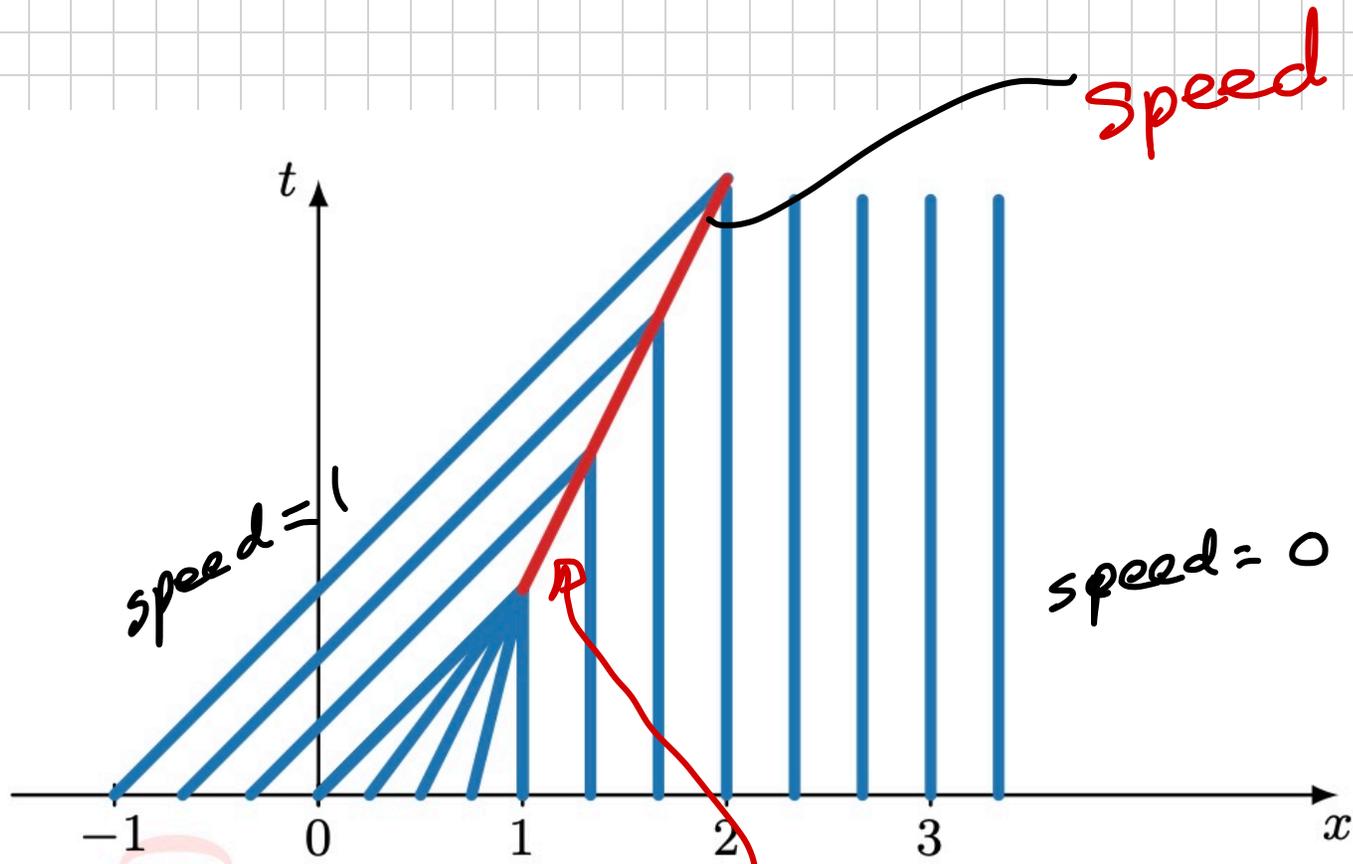
Cases

① Shock  $u_0(x, 0) = \begin{cases} 1 & x < 0 \\ 1-x & x \in [0, 1] \\ 0 & x > 1 \end{cases}$

② Rarefaction

$$u_0(x, 0) = \begin{cases} 0 & x < 0 \\ x & x \in [0, 1] \\ 1 & x > 1 \end{cases}$$

Draw the characteristic curves.



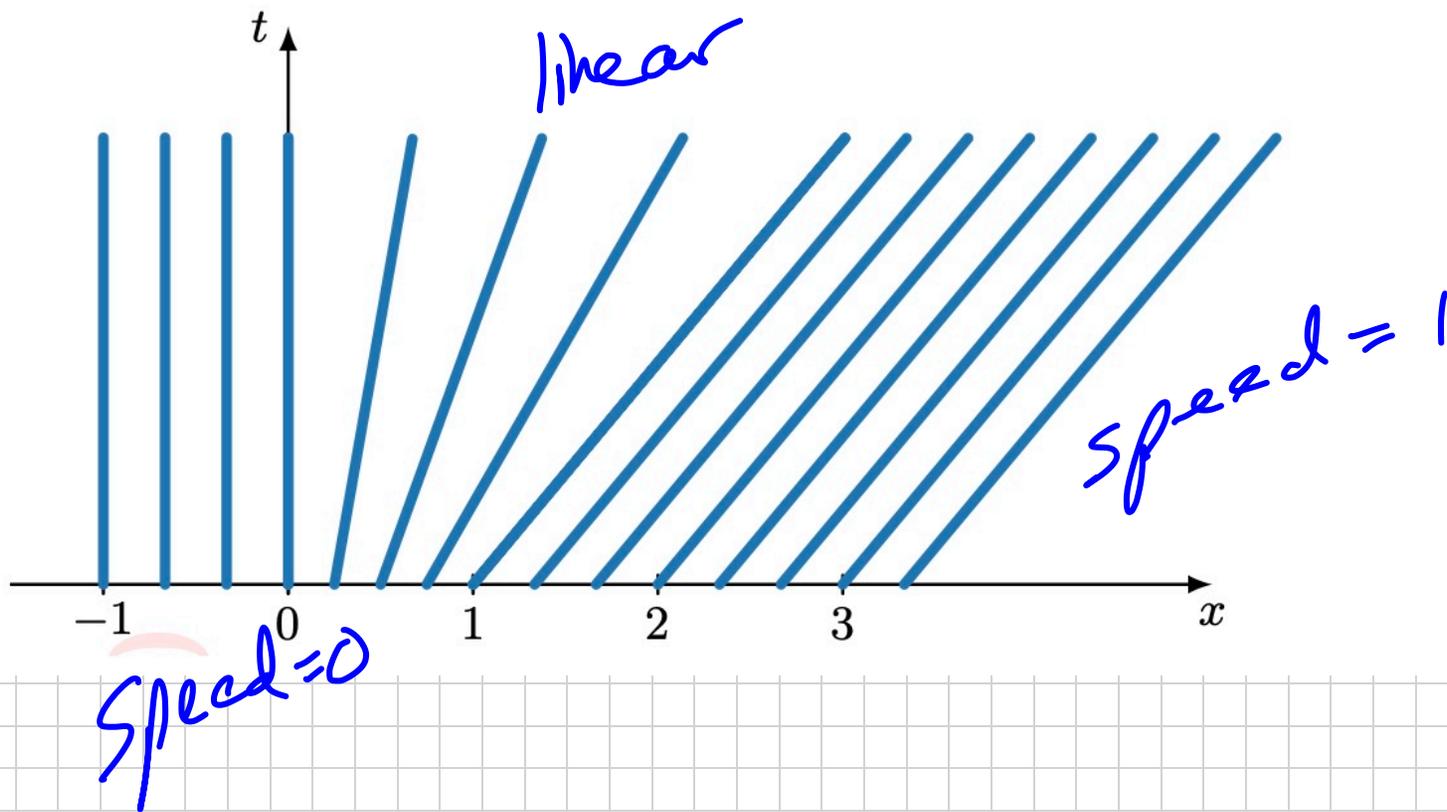
let  $\hat{x}(t) = \text{curve}$

then  $\hat{x}'(t) = \frac{f(u^+) - f(u^-)}{u^+ - u^-}$

on the right
on the left

$$= \frac{\frac{1}{2} \cdot 0^2 - \frac{1}{2} \cdot 1^2}{0 - 1} = \frac{1}{2}$$

Burgers:  $f(u) = \frac{u^2}{2}$

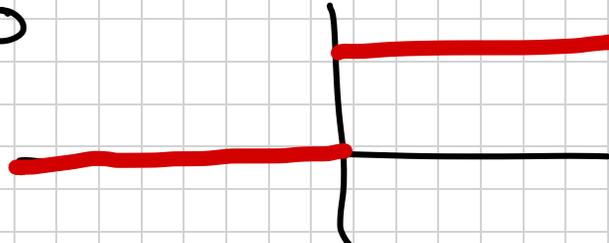


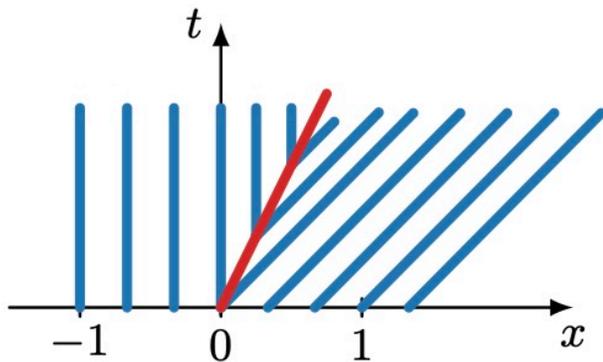
### Definition 6.5: Riemann Problem

Consider scalar conservation law (6.13) with initial condition

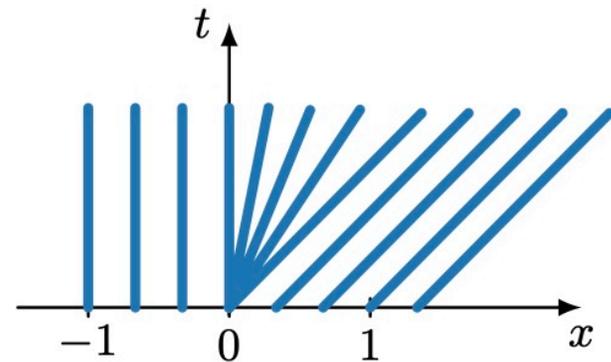
$$u(x, 0) = \begin{cases} u^- & \text{if } x \leq 0, \\ u^+ & \text{if } x > 0. \end{cases} \quad (6.29)$$

This problem consisting of constant left and right states separated by a jump discontinuity is called a Riemann problem.

$$u(x, 0) = \begin{cases} 0 & x \leq 0 \\ 1 & x > 0 \end{cases}$$




a. Unphysical weak solution with entropy-violating shock.



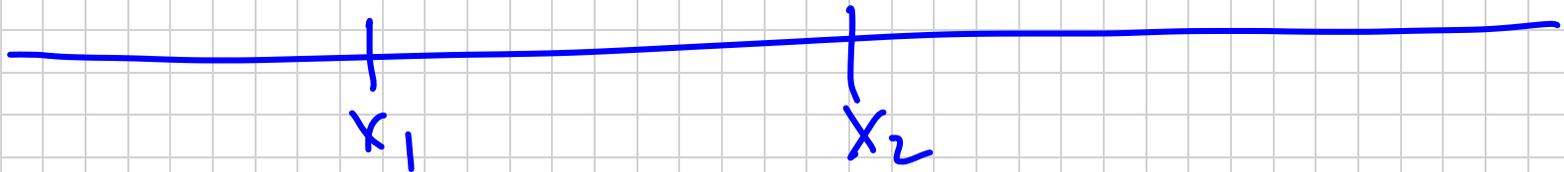
b. Unique physically relevant, vanishing viscosity weak solution (rarefaction wave).

Return to conservation law:

Conserved quantity.

$u(x, t)$  = density at  $x$   
at time  $t$ .  $\left[\frac{\text{kg}}{\text{m}}\right]$

$$\int_x^{x_2} u(x, t) dx = \text{total mass } M(x_1, x_2) \text{ at time } t.$$



mass can only change due to  
flow in or out at  $x_1$  or  $x_2$   
"flux"

$$F_1(t) = \text{rate at } x_1 \quad \left[\frac{\text{kg}}{\text{s}}\right]$$

$$F_2(t) = \text{rate at } x_2$$

(flux "to the right"  $F_i(t) > 0$ )

$$\frac{d}{dt} \int_{x_1}^{x_2} u(x,t) dx = \underbrace{F_1(t)}_{\text{flow in at } x_1} - \underbrace{F_2(t)}_{\text{flow out at } x_2}$$

general

$$= f(u(x_1, t)) - f(u(x_2, t))$$

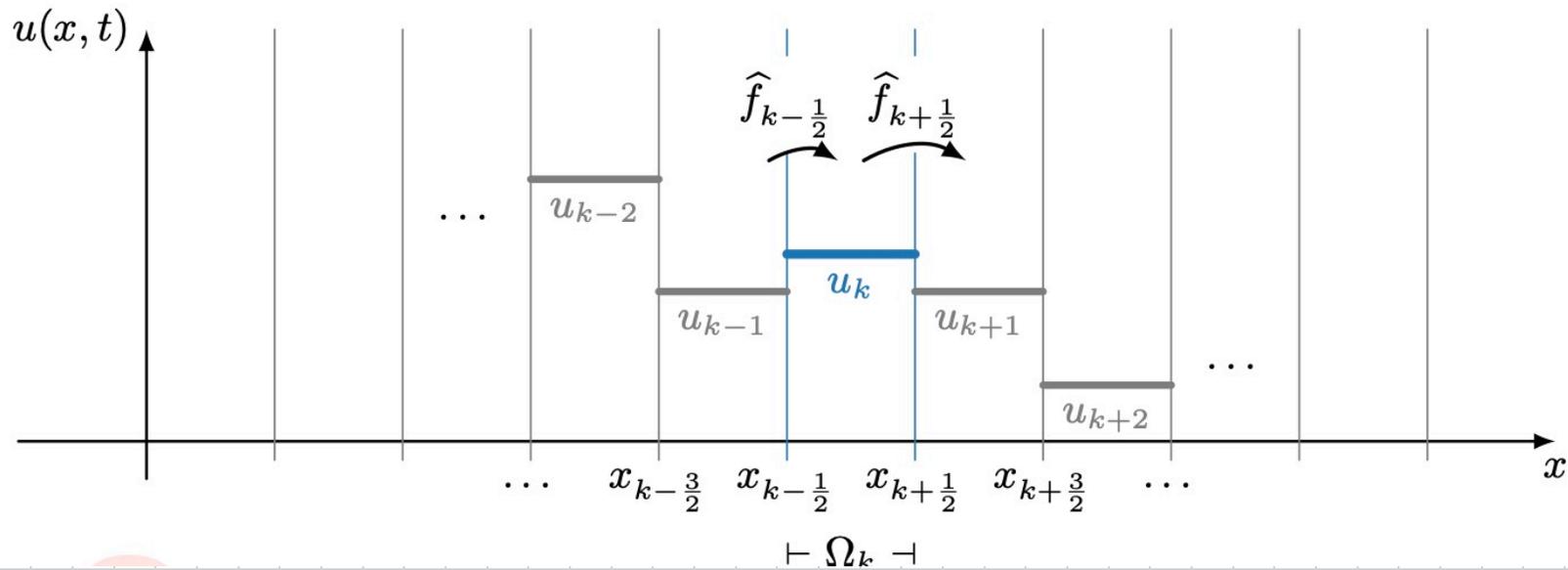
$$= -f(u(x, t)) \Big|_{x_1}^{x_2}$$

(Assume smooth)

$$\Rightarrow \frac{d}{dt} \int_{x_1}^{x_2} u(x,t) dx + \int_{x_1}^{x_2} \frac{d}{dx} f(u(x,t)) dx = 0$$

$$\Rightarrow \int_{x_1}^{x_2} u_t + (f(u))_x dx = 0$$

Idea



$$\Omega_k = [x_{k-1/2}, x_{k+1/2}]$$

$$\int_{\Omega_k} \frac{\partial u}{\partial t} + \frac{\partial f(u)}{\partial x} dx = 0$$

$$h_x \frac{d \bar{u}_k(t)}{dt} + f(u(x_{k+1/2}, t)) - f(u(x_{k-1/2}, t)) = 0$$

$$\bar{u}_k(t) = \frac{1}{h_x} \int_{\Omega_k} u(x, t) dx$$

The solution  $u(x,t)$  in  $\Omega_k$  satisfies

$$\frac{d \bar{u}_k(t)}{dt} + \frac{f(u(x_{k+1/2}, t)) - f(u(x_{k-1/2}, t))}{hx} = 0$$

exactly.

To turn into a method:

① discrete values  $u_k \approx \bar{u}_k(t)$

② define a numerical flux:

$$\hat{f}_{k+1/2} \approx f(u(x_{k+1/2}, t))$$

③ ODE solver

method "framework":

$$\frac{d u_k(t)}{dt} + \frac{\hat{f}_{k+1/2}(t) - \hat{f}_{k-1/2}(t)}{h_x} = 0$$



ODE solver

FWD Euler

RK4

whatever

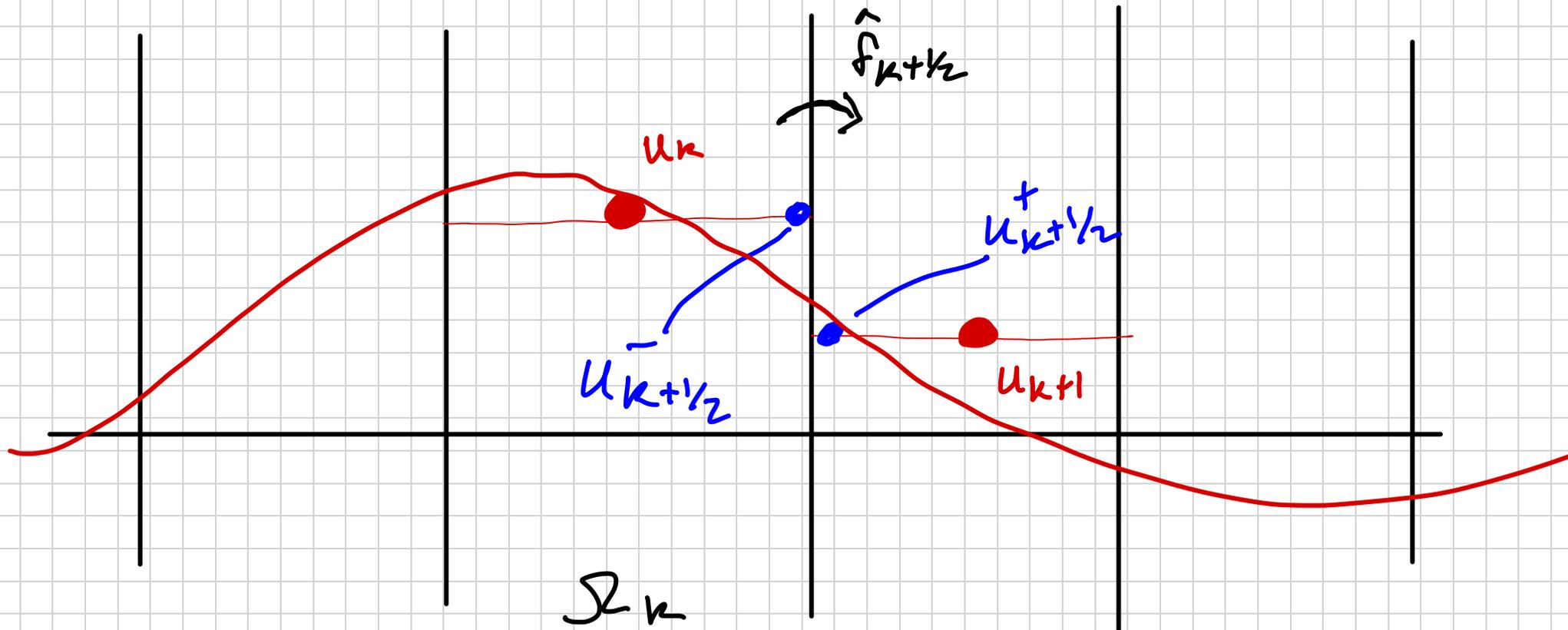
For  $\hat{f}_{k+1/2}(t)$ .

We will use

$$\hat{f}_{k+1/2}(t) \approx f^*(\cdot, \cdot)$$

$$= f^*(u_{k+1/2}^-(t), u_{k+1/2}^+(t))$$

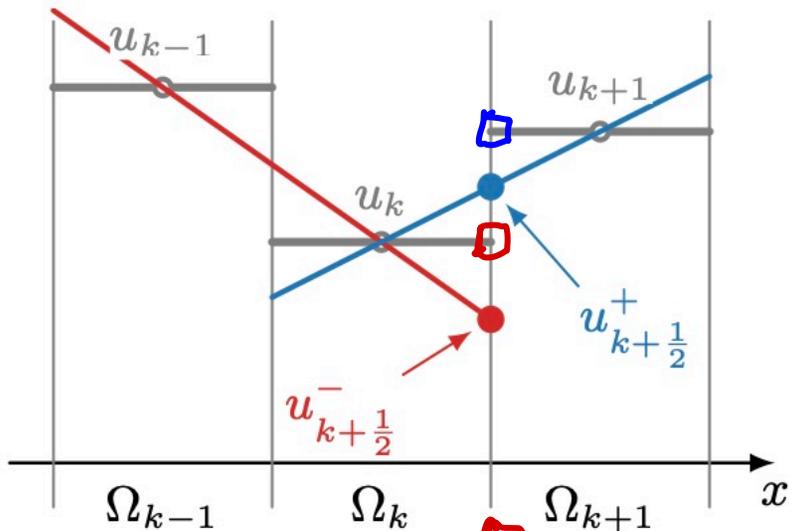
approx value "just to the left" of  $x_{k+1/2}$



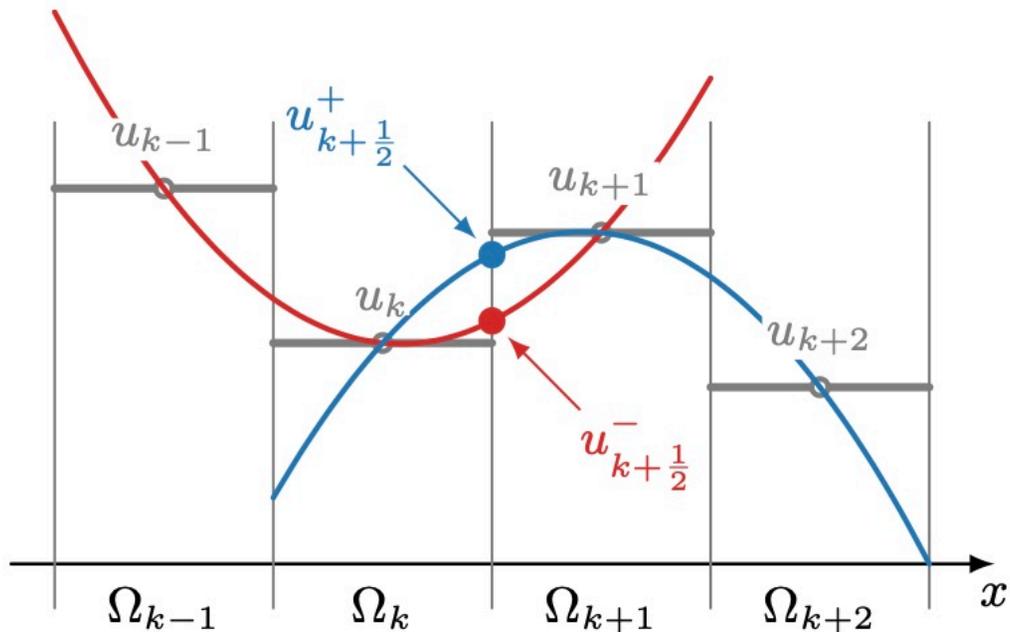
easiest:

$$u_{k+1/2}^- = u_k$$

$$u_{k+1/2}^+ = u_{k+1}$$



$x_{k+1/2}$



let  $u_{k,r} =$  approx cell average at  $t$   
+  
FWD Euler

$$\frac{u_{k,r+1} - u_{k,r}}{\Delta t} + \frac{f^*(u_{k,r}, u_{k+1,r}) - f^*(u_{k-1,r}, u_{k,r})}{\Delta x} = 0$$

The PDE gives us  $f(u)$

### Definition 6.8: Consistent numerical flux function

A numerical flux function  $f^*(u^-, u^+)$  for conservation law (6.40) is called consistent if

- $f^*(u, u) = f(u) \quad \forall u$
- $f^*(u^-, u^+)$  is Lipschitz continuous in each argument.

advection

$$u_t + (au)_x = 0$$
$$f(u) = a \cdot u$$

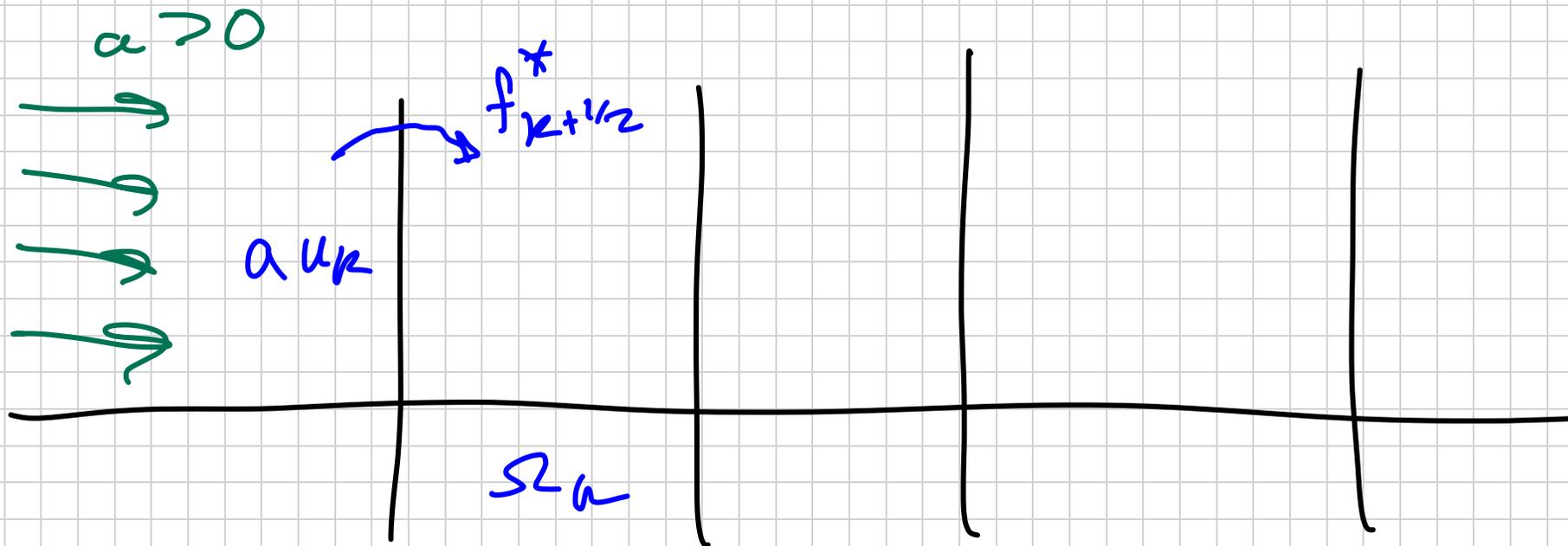
Pick  $f^*(u_{k-1}, u_{k+1}) = a u_k$

and  $f^*(u_{k-1}, u_k) = a u_{k-1}$

$$0 = \frac{u_{k+1} - u_k}{h_t} + \frac{f^*(u_k, u_{k+1}) - f^*(u_{k-1}, u_k)}{h_x}$$

$$\Rightarrow \frac{u_{k+1} - u_k}{h_t} + a \frac{u_k - u_{k-1}}{h_x} = 0$$

assume  $a > 0$



$$\text{if } a > 0 : f^*(u_{k2}, u_{k+12}) = a u_{k2}$$

$$\text{if } a < 0 : f^*(u_{k2}, u_{k+12}) = a u_{k+12}$$

Q: what if we want  $f^*$  to "work" for both?

hint: use  $|a|$

$$\rightarrow f^*(u_{k2}, u_{k+12}) = a \left( \frac{u_{k2} + u_{k+12}}{2} \right) - \frac{|a|}{2} (u_{k+12} - u_{k2})$$